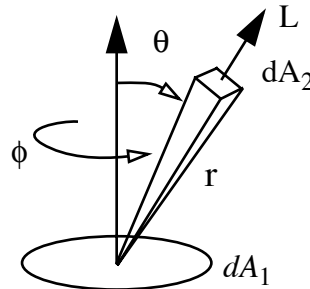


# RADIOMETRY OF LAMBERTIAN SOURCES

## Relate $M_\lambda$ and $L_\lambda$

- Valuable exercise in spherical coordinate integration



- Area element of sphere, radius  $r$

$$\begin{aligned} dA_2 &= r d\phi r \sin\theta d\theta \\ &= r^2 \sin\theta d\theta d\phi \end{aligned}$$

- Start with radiance to get flux within area  $dA_2$

$$L(\theta, \phi) = \frac{d^2\Phi}{dA_1 \cos\theta d\Omega}$$

$$d^2\Phi = L(\theta, \phi) \cos\theta dA_1 d\Omega$$

$$= L(\theta, \phi) \cos\theta dA_1 \frac{dA_2}{r^2}$$

$$= L(\theta, \phi) \cos\theta dA_1 \sin\theta d\theta d\phi$$

- **Total flux in hemisphere above source**

$$\Phi = dA_1 \int_0^{2\pi} d\phi \int_0^{\pi/2} L(\theta, \phi) \cos\theta \sin\theta d\theta$$

*Now, in general  $L = L(\theta, \phi)$ , but for a Lambertian source*

***L = constant***

- ***Therefore***

$$\begin{aligned}\Phi &= dA_1 2\pi L \frac{(\sin\theta)^2}{2} \Big|_0^{\pi/2} \\ &= \pi L dA_1\end{aligned}$$

***Radiant exitance definition***

$$\begin{aligned}M &= \frac{\Phi}{dA_1} \\ &= \pi L\end{aligned}$$

## *Include spectral variation*

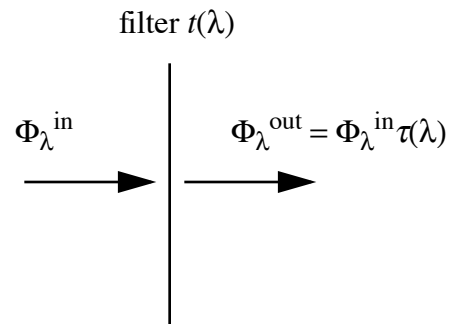
$$M_\lambda = \pi L_\lambda$$

$$\text{or } L_\lambda = M_\lambda / \pi$$

- *Can think of  $\pi$  as having units of  $sr$ , which cancels the  $sr^{-1}$  units of  $L$*
- *Note:  $M \neq 2\pi L$ , as one might guess since hemisphere is  $2\pi sr$*   
*Why?*

# TRANSMITTANCE

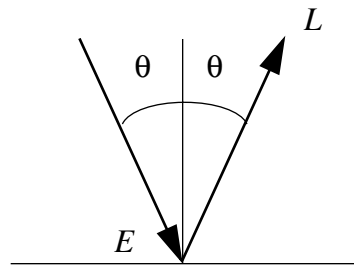
*Defined as ratio of transmitted output flux to input flux*



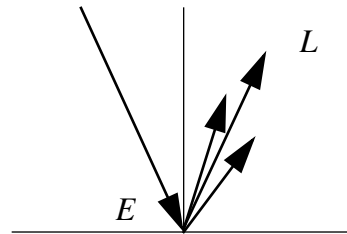
$$t(\lambda) = \Phi_{\lambda}^{\text{out}} / \Phi_{\lambda}^{\text{in}}$$

# REFLECTANCE

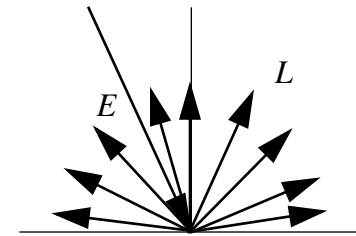
*Three types:*



*specular (mirror)*



*directional*



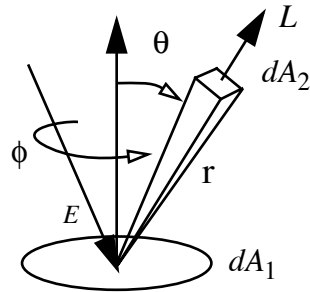
*diffuse (Lambertian)*

*Most natural surfaces are approximately Lambertian*

*for  $\theta < 40^\circ$ , snow and sand for  $\theta < 60^\circ$*

- *At larger  $\theta$ , natural surfaces tend to become directional*

## *Reflectance of Lambertian surface*



- *Reflectance defined similarly to transmittance*

$$\rho(\lambda) = \Phi_{\lambda}^{out} / \Phi_{\lambda}^{in}$$

- *From earlier derivation*

$$\Phi_{\lambda}^{out} = \pi L_{\lambda} dA_1$$

*and*

$$\Phi_{\lambda}^{in} = E_{\lambda} dA_1$$

*Therefore,*

$$\rho(\lambda) = \pi L_{\lambda} / E_{\lambda}, \quad 0 \leq \rho(\lambda) \leq 1$$

*Given reflectance of surface and incident irradiance,  
can get radiance*

$$L_{\lambda} = \rho(\lambda) E_{\lambda} / \pi$$



## BAND-AVERAGED IRRADIANCE

*Cascade spectral quantities from source-to-surface-to-surface . . .*

$$E_{\lambda} = (\text{geometric factors}) \cdot L_{\lambda} \cdot t_1(\lambda) \cdot t_2(\lambda) \dots \rho_1(\lambda) \dots$$

- *until arriving at a spectrally-integrating element, i.e. a detector,*

$$E_{total} = \int E_{\lambda} \cdot S(\lambda) d\lambda$$

*where  $S(\lambda)$  is the detector's **spectral sensitivity***

*Example of a specific detector  $S(\lambda)$  is human vision system sensitivity  $V(\lambda)$*

*$E_{total}$  is the **effective** irradiance, because it is what is*

*detected and measured*

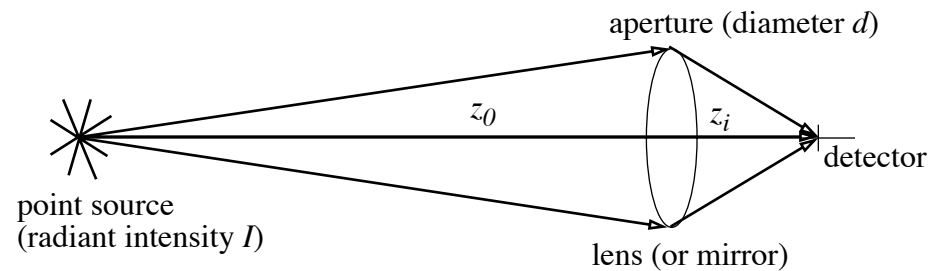
*Band-averaged spectral irradiance*

$$E_b = \frac{E_{total}}{\int_0^{\infty} S(\lambda) d\lambda}$$

# RADIOMETRY OF OPTICAL SYSTEMS

## telescope

- *collects light from point source*
- *“light bucket”*



- *assume:*
  - *source and detector on optical axis*
  - *object-to-sensor distance much greater than aperture*

**diameter**  $z_0 \gg d$

- **no transmission losses**
- **irradiance at lens aperture**  $E = I/z_0^2$  (**inverse square law**)
- **flux collected by aperture**

$$\begin{aligned}\Phi_c &= E \cdot A_{\text{aperture}} \\ &= E \cdot \frac{\pi d^2}{4} \\ &= \frac{I}{z_0} \cdot \frac{\pi d^2}{4}\end{aligned}$$

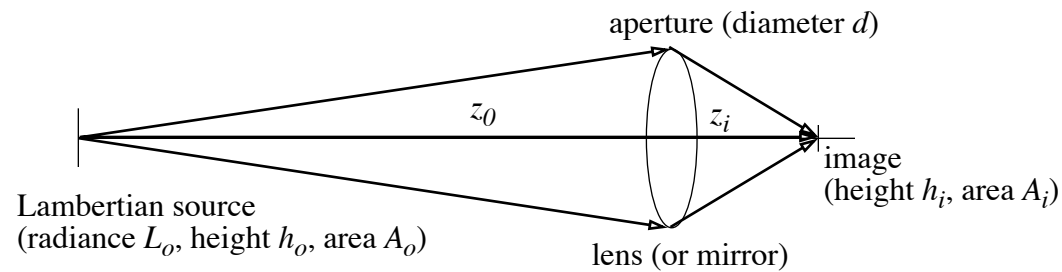
- **flux at detector (no losses)**  $\Phi_i = \Phi_c$ 
  - **proportional to radiant intensity of source**
  - **proportional to square of aperture diameter**

- *inversely proportional to square of object-to-sensor distance*

*What is the only way to increase the amount of light collected by a telescope?*

## *Camera*

- *assume:*
  - *Lambertian source and detector plane normal to optical axis*
  - *magnification*  $m = h_i/h_o = z_i/z_o$



- **flux collected by aperture**

$$\begin{aligned}\Phi_c &= L_o A_o \Omega \\ &= L_o A_o \frac{A}{z_o^2} \\ &= \frac{L_o A_o}{z_o^2} \cdot \frac{\pi d^2}{4}\end{aligned}$$

- **NOTE: similarity to “light bucket” equation**
- **flux at detector (no losses)**  $\Phi_i = \Phi_c$

- **irradiance at detector**

$$\begin{aligned}
 E_i &= \Phi_c / A_i \\
 &= \frac{\Phi_c}{m^2 A_o} \\
 &= L_o \frac{\pi d^2}{m^2 4z_o^2}
 \end{aligned}$$

- **by definition of magnification (see Geometrical Optics)**

$$E_i = L_o \frac{\pi d^2}{4z_i^2}$$

- **define *effective f-number N* of sensor,**  $N = z_i/d$

- **then,**

$$E_i = L_o \frac{\pi}{4N^2} \quad \text{Camera Equation}$$

- **NOTE:**

- **proportional to radiance of source**

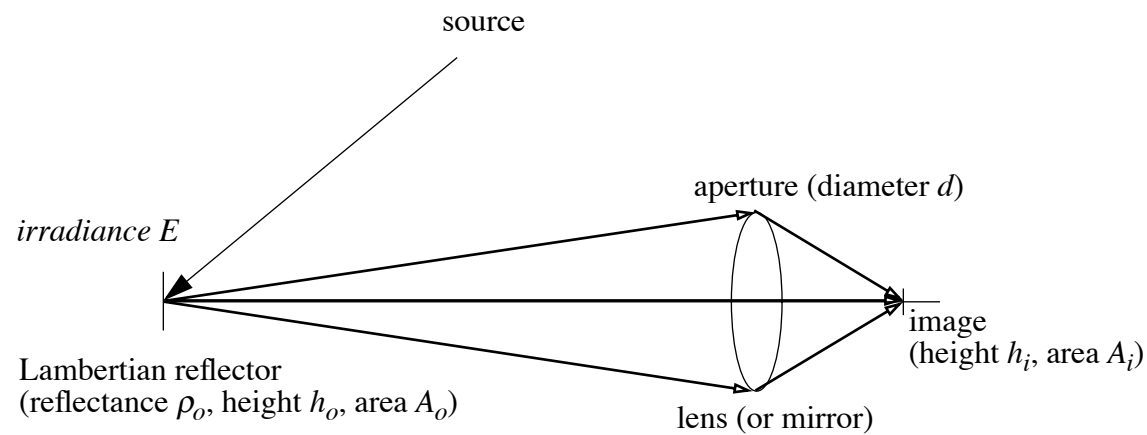
- *inversely proportional to square of sensor f-number*
- *does not depend on  $z_o$ , the source-to-sensor distance*
- ***f-number  $N$  typically preset on camera to 1.4, 2, 2.8, 4, 5.6, 8, 11, 16, 22***

*What is rationale for the above preset values of  $N$ ?*



## Camera imaging *reflecting object*

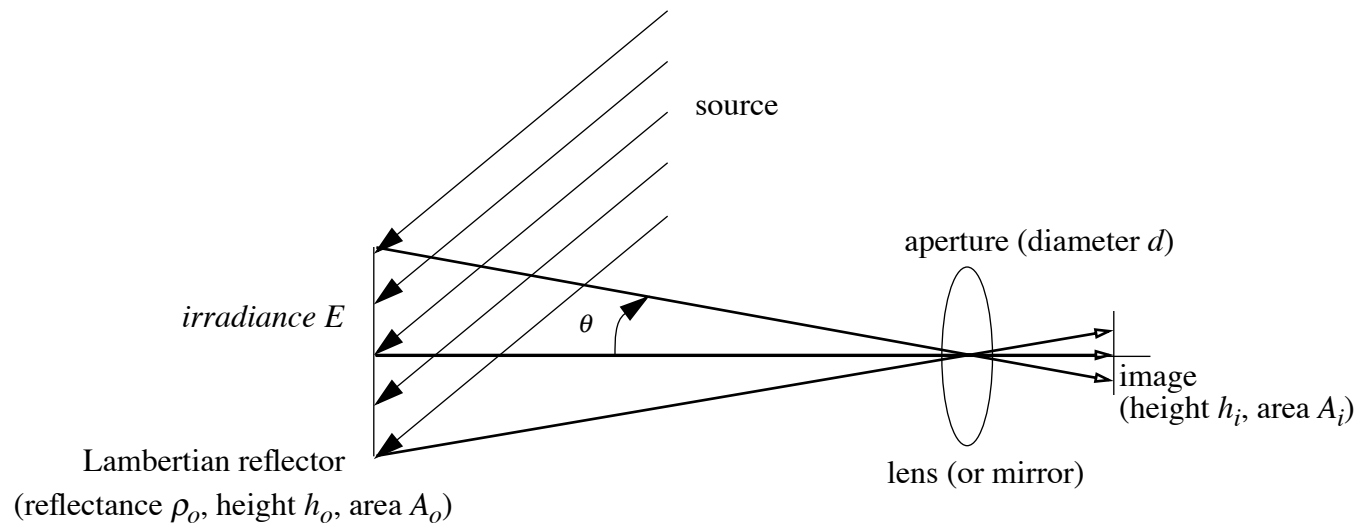
- **Irradiance on object  $E$**
- **Reflectance of object  $\rho_o$**



- **From earlier derivation:**

$$E_i = \frac{E_o \rho}{4N^2}$$

## Camera imaging *extended reflecting object*



- **$\text{Cos}^4\theta$  Law applies:**

$$E_i(\theta) = \frac{E_o \rho}{4N^2} (\cos\theta)^4$$

**With object irradiance and reflectance fixed, what is the only way to increase the**

*amount of light collected by a camera?*